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Structural, Electrical and Piezoelectric Properties of $Ba_{1-x}Nd_xTiO_3$ Ceramic for Microelectronic Applications

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ABSTRACT

Neodymium (Nd)-doped barium titanate (BaTiO₃) ceramics with general formula Ba_{1-x}Nd_xTiO₃ ($0 \le x \le 0.13$) were synthesized by the conventional solid-state reaction method at 1400 °C for 36 hours in air. A single-phase structure of Nd-doped BaTiO₃ was observed from $0 \le x \le 0.10$. The tetragonal distortion of undoped BaTiO₃ progressively decreased with increasing Nd content (x). The samples retain a tetragonal structure for $x \le 0.015$, transformed to a cubic phase for $0.03 \le x \le 0.1$ and showed to coalesce at x = 0.13. The undoped BaTiO₃ exhibited a maximum permittivity approximately $\varepsilon = 8500$ at the Curie temperature (T_c), while a slightly reduced value of about $\varepsilon = 8200$ for Nd at x = 0.005. Nd substitution resulted in a shift of the permittivity maximum towards lower T_c within the tetragonal region, whereas compositions in the cubic displayed a nearly linear temperature dependence without a distinct maximum. The composition with x = 0.005 showed the highest piezoelectric charge coefficient, $d_{33} = 49$ pC/N and piezoelectric voltage coefficient of $g_{33} = 0.2 \times 10^{-2}$ V.m/N.

Keywords: Neodymium doped barium titanate, Permittivity, Rietveld refinement, BaTiO₃

1. INTRODUCTION

Researchers are attempting to reduce the size of all electronic devices to be as small and lightweight as possible. Due to this trend, high-permittivity materials such as barium titanate (BaTiO₃) have become increasingly important in ceramic materials. BaTiO₃, with its outstanding dielectric properties and high thermal stability, is widely used in the fabrication of electronic devices such as multilayer ceramic capacitors (MLCCs), positive temperature coefficient of resistance (PTCR) thermistors, piezoelectric sensors, transducers, and more [1–2].

BaTiO $_3$ exhibits a perovskite structure with the general formula ABO $_3$, can have its properties tailored through doping at either the A-site or B-site. Doping in BaTiO $_3$ is strongly influenced by the ionic radii of the dopant elements [3]. Recently, trivalent rare-earth elements have been reported as effective dopants to improve the physical and electrical properties of BaTiO $_3$ [4,5]. Morrison et al. [4], in his study, discovered an unusually high permittivity of approximately 25,000 in BaTiO $_3$ when doped with lanthanum (La) at both sites at high frequencies. This exceptionally high permittivity was accompanied by a trend of shifting the Curie temperature (T_c) to lower values. Due to its larger ionic radius, La appears to substitute exclusively at the A-site.

Neodymium (Nd) is another promising trivalent rare-earth element used as a dopant. The Nd $^{3+}$ ion (1.27 Å) is presumed to occupy the Ba $^{2+}$ site (1.61 Å) rather than the Ti $^{4+}$ site (0.605 Å), due to size incompatibility. However, double substitution of Nd-doped BaTiO $_3$, as reported by Hirose et al. [5], resulted in structural changes from the tetragonal to cubic phase with increasing Nd content. In addition, the sharp maximum in permittivity broadens rapidly with increasing Nd and gradually shifts to lower temperatures. Nevertheless, there is limited information on the variation of lattice parameters, unit cell volume, complete structural refinement, and detailed electrical and piezoelectric properties of Nd-doped BaTiO $_3$.

Therefore, in this paper, a systematic study on the effect of Nd doping at the A-site in $BaTiO_3$ is presented. Rietveld refinement, as well as electrical and piezoelectric characterizations, are reported. The aim is to examine and, where necessary, optimize the electrical properties of sintered $BaTiO_3$ and Nd-doped $BaTiO_3$ for their potential in piezo electronic applications.

2. EXPERIMENTAL PROCEDURE

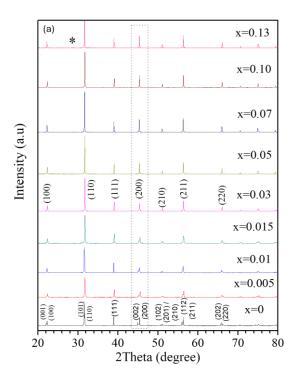
Ba_{1-x}Nd_xTiO₃ ceramics were synthesized via the conventional solid-state reaction method from high-purity commercial precursors: barium carbonate (BaCO₃) (99%), titanium dioxide (TiO2) (99%), and neodymium (III) oxide (99%)obtained (Nd₂O₃)from Sigma-Aldrich. Stoichiometric amounts of the powders were weighed, thoroughly mixed in an agate mortar and pestle to ensure chemical homogeneity and then calcined in air at 1000 °C for 12 hours. The calcined powders were reground and uniaxially pressed into disc-shaped pellets (13 mm diameter, ~1 mm thickness) under a pressure of 5 tons. These pellets were sintered at 1400 °C for 6 hours in a programmable muffle furnace to achieve densification. Phase identification and structural analysis were performed using a Bruker D2 Phaser X-ray diffractometer equipped with a LYNXEYE 1D detector and Cu-Kα radiation. Data were collected at room temperature using a step size of 0.02° and a dwell time of 0.2 seconds. Rietveld refinements were conducted using the GSAS/EXPGUI software package[6-7], and crystallographic models were visualized with VESTA software [8], following refinement strategies from established literature [9-10]. The sintered pellets were polished and coated with silver paste to form electrodes, then fired to ensure proper electrical contact. Electrical property measurements, including dielectric properties, capacitance, and resistance, were carried out using an IM3570 HIOKI Impedance Analyzer over a frequency range of 10 Hz to 1 MHz and a temperature range of 40 °C to 200 °C, using a programmable furnace with a controlled heating rate of 5 °C/min.

3. RESULTS AND DISCUSSION

3.1. X-ray Diffraction Analysis

Figure 1(a) presents the X-ray diffraction (XRD) patterns of 0.13). All compositions up to x = 0.10 exhibit a single-phase perovskite structure without detectable secondary phases, indicating successful incorporation of Nd3+ into the BaTiO3 lattice. The diffraction pattern of undoped BaTiO₃ corresponds to a tetragonal structure with refined lattice parameters a = 3.9937 Å and c = 4.0342 Å, in close agreement with previously reported values (a = 3.994 Å, c = 4.038 Å) [11-12]. With Nd doping in the range of $0.03 \le x \le 10^{-10}$ 0.10, a structural transition from tetragonal to pseudo-cubic symmetry become observed. The characteristic peak splitting of the (002) and (200) reflections, indicative of tetragonality, gradually merges into a single peak as shown in the enlarged view in Figure 1(b). This behaviour suggests a progressive suppression of tetragonal distortion with increasing Nd concentration [13].

At low doping levels (0 < x < 0.015), a decrease in tetragonality is observed. The composition at x = 0.015 exhibits cubic symmetry with a lattice parameter of a = 4.0014 Å. As Nd content increases up to x = 0.10, the cubic structure is maintained and the lattice parameter increases linearly to 4.0050 Å. For x \geq 0.13, secondary phases begin to emerge, with additional diffraction peaks attributed to Nd₂Ti₂O₇ [13]. This indicates that the solubility limit of Nd in the BaTiO₃ lattice lies between x = 0.1 and 0.13. At x = 0.13, the primary phase remains cubic with a = 3.9995 Å, although the presence of impurity peaks confirms phase segregation beyond the solid solution limit.



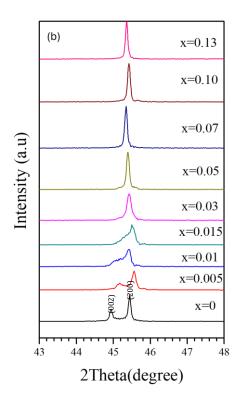


Figure 1. (a) XRD patterns of $Ba_{1-x}Nd_xTiO_3$ ceramics (b) Enlarge Xray Splitting XRD peak splitting for (002) and (200) reflection.

The evolution of lattice parameters as a function of Nd content is shown in Figure 2. For $0 \le x \le 0.015$, the lattice parameter an increase while c decreases slightly, resulting in a reduction in the c/a ratio and a corresponding decrease in unit cell volume. For $0.03 \le x \le 0.10$, the lattice parameter a remains nearly constant, consistent with a stable cubic

phase. These variations are attributed to the substitution of smaller Nd^{3+} ions (1.27 Å) for larger Ba^{2+} ions (1.61 Å) at the A-site, leading to lattice contraction. The deviation from linearity in lattice parameters observed at $x \ge 0.13$ further supports the existence of a solubility limit. These results are consistent with those reported by Yao et al. [14].

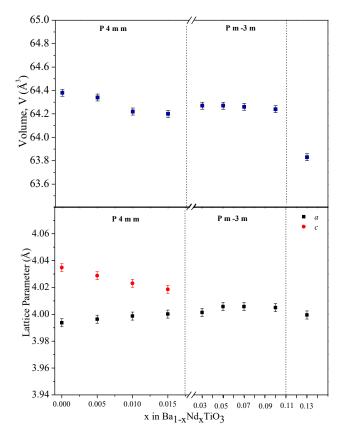


Figure 2. Lattice parameter and unit cell volume of $Ba_{1-x}Nd_xTiO_3$ ($0 \le x \le 0.13$).

In the refinement analysis, common parameters that can be refined include the scale factor, background, lattice parameters, 2θ zero, profile parameters, atomic positions, and thermal parameters, as reported elsewhere [9]. Table 1 presents the refinement data for tetragonal Nd-doped BaTiO₃. Small variations in the atomic positions of titanium and oxygen were observed, while the positions of barium and neodymium remained fixed. The substitution of Nd³⁺ creates Ba-site deficiencies to maintain charge neutrality, which increase with higher Nd content. Consequently, the tetragonal unit cell undergoes shrinkage, displacing the Ti⁴⁺ ions from the centre of the octahedral sites and weakening

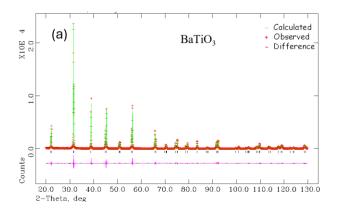
the TiO_6 octahedra. This structural distortion initiates a phase transition from tetragonal to cubic. Although the R-factors and χ^2 values are relatively higher than those of undoped BaTiO3, they remain within acceptable limits. Table 2 shows the R-factors (R_{wp} and R_p) for cubic Nd-doped BaTiO3, which are considered acceptable and satisfactory (below 20%). Moreover, the χ^2 values are relatively low and below 5, indicating good fit quality [15]. The R_p and R_{wp} values confirm a reasonable agreement between the refined and observed XRD patterns for all $Ba_{1-x}Nd_xTiO_3$ ceramics, in both tetragonal and cubic phases. These results suggest that the structural refinement is reliable.

Table 1 Structural refinement data for Ba1-xNdxTiO3 ($0 \le x \le 0.015$)

Doping Concentration	x = 0.005	x = 0.01	x = 0.015
Ti, z	0.5169(8)	0.5142(1)	0.5136(1)
Ba (1a), U _{iso}	0.0076(3)	0.0058(3)	0.0069(3)
Nd (1a),U _{iso}	0.0139(2)	0.0169(2)	0.0059(2)
Ti (1b), U _{iso}	0.0080(1)	0.0069(1)	0.0057(1)
01(1b), U _{iso}	0.0120(2)	0.0175(2)	0.0070(2)
02(2c),U _{iso}	0.0010(2)	0.0095(4)	0.0142(2)
Rwp (%)	13.89	14.5	13.6
R _p (%)	11.04	13.1	10.5
χ^2	4.47	4.93	4.51

Table 2 Structural refinement data $Ba_{1-x}Nd_xTiO_3$ (0.03 $\leq x \leq 0.10$)

Doping Concentration	x = 0.05	x = 0.07	x = 0.10
Ti, z	-	-	-
Ba (1a), U _{iso}	0.0091(2)	0.007(2)	0.0072(2)
Nd (1a),U _{iso}	0.0075(2)	0.0069(2)	0.0174(2)
Ti (1b), U _{iso}	0.0079(2)	0.0051(2)	0.0054(2)
01(1b), U _{iso}	0.0131(2)	0.0103(1)	0.0128(2)
02(2c),U _{iso}	-	-	-
Rwp (%)	11.9	9.34	8.58
R _p (%)	9.04	7.24	4.99
χ^2	3.33	2.43	2.22



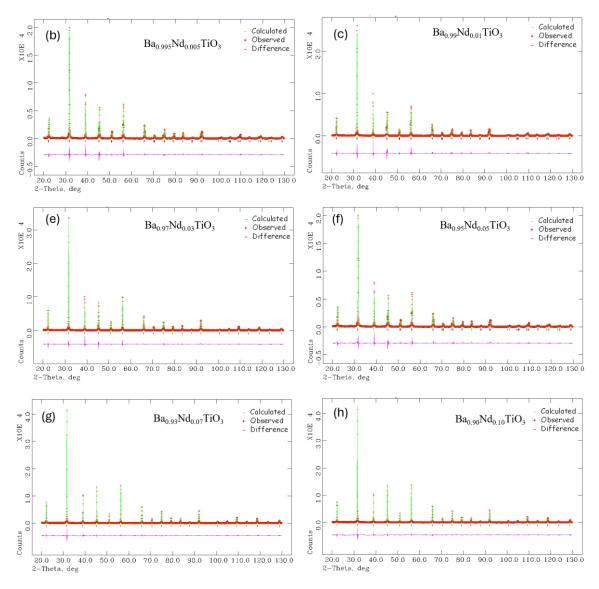


Figure 3. Plot of structural refinement of Ba1-xNdxTiO3 [a-h] ($0 \le x \le 0.10$).

The structural plot of Nd-doped $BaTiO_3$ in Figure 3 demonstrates good agreement between the experimental XRD patterns and the refined model.

3.2. Impedance Spectroscopy Analysis

Figure 4 (a-c) shows the variation in the maximum permittivity of Nd-doped BaTiO₃ ceramics as a function of temperature. Significant changes were observed upon addition of Nd³⁺. The maximum permittivity sharp peak was lowered down with increasing Nd³⁺. The decrease in permittivity was accompanied by a shift of T_c to lower temperature. The variation of permittivity with temperature shows a broadened maximum peak with increasing Nd³⁺ concentration. Undoped BaTiO₃ exhibited

the highest permittivity, reaching approximately $\epsilon \approx 8500$ at the Curie temperature (T_c) of 120 °C. However, with Nd doping (0.005 \leq x \leq 0.015), the maximum permittivity decreased to approximately 8200, 5800, and 2300, respectively, as measured at 1 kHz. The results show that the Nd³+ addition shifted the temperature of the maximal permittivity from 120 °C (x = 0) below than that for 0.03 < x < 0.10. In contrast, samples with Nd content in the range 0.03 < x < 0.10 exhibited nearly constant permittivity over the measured temperature range. The lower permittivity in the cubic region may be attributed to its location in the paraelectric region, where the material retains no residual polarization [16].

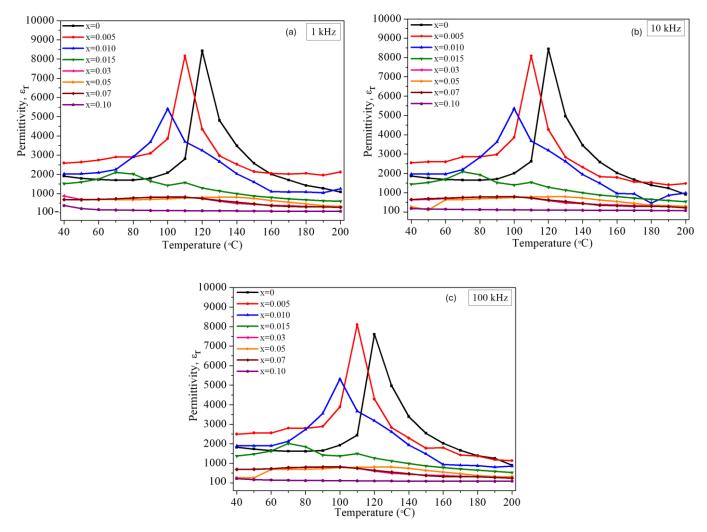


Figure 4. Permittivity of Ba_{1-x}Nd_xTiO₃ ($0 \le x \le 0.13$) at (a) 1 kHz (b) 10 kHz and (c) 100 kHz.

There was a decrease in permittivity with increasing frequency as shown in Figure 4. This trend also was previously observed by Ramirez [16], Yao et al., [14], Ganguly et al., [17], and Sharma et al. [18]. Permittivity value was slightly higher at low frequencies than at higher frequencies. At low frequencies, space charge polarization may become dominant contribution samples exhibiting very high permittivity, which is strongly related to dipole orientation, is unable to keep pace with rapidly varying electric field frequency. Thus, permittivity begins to decrease with increasing frequency due to the phase lag between dipole alignment and electric field [19].

Figure 5 shows the temperature dependence in dielectric loss of Nd-doped $BaTiO_3$ at three different frequencies. The dielectric loss of the $Ba_{1-x}Nd_xTiO_3$ is very low at much lower temperature remaining below 0.8 at 1 kHz. The dielectric loss is very low below Tc and increase slightly near Tc. However, the dielectric loss becomes very minimal as measured at frequency 10 kHz and 100 kHz. Dielectric loss of any material is related to the heat dissipated by the material, when electrical field is applied [20]. By doping with Nd, the dielectric loss of $BaTiO_3$ markedly drop to a lower value, from 0.10 to 0.004 at room temperature. This means that lower dielectric loss is preferable especially for device performance and efficiency.

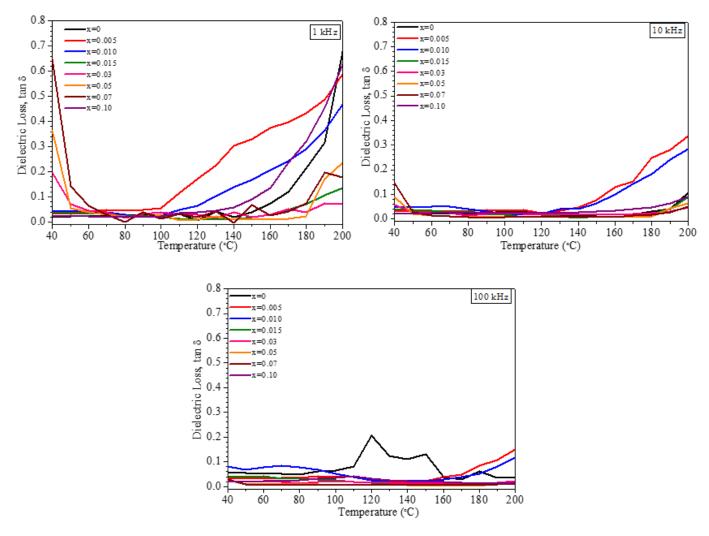


Figure 5. Dielectric loss of Ba1-xNdxTiO3 ($0 \le x \le 0.13$) at 1 kHz, 10 kHz and 100 kHz.

Figure 6(a) shows the frequency dependence conductivity of tetragonal Nd-doped BaTiO3 and cubic Nd-doped BaTiO3 in Figure 6(b) across frequency of 1 kHz to 100 kHz. Log σ versus log frequency below 1 kHz were excluded, as it essentially shows noisy data since sample impedance lies outside the measuring range of instrumentation [21]. For pure BaTiO3, the linear region of log σ at 10 Hz to 1 kHz belonged to DC conductivity (which was excluded from extracted frequency due to noisy data. Above 1 kHz, log σ was the essence of the AC conductivity. The conductivity of

BaTiO₃ increased linearly with the frequencies from 1 kHz to 100 kHz. Overall conductivity of tetragonal Nd-doped BaTiO₃ lies in the range of 10^{-5} to 10^{-8} Scm⁻¹. The addition of x seemed to influence the original conductivity of pure BaTiO₃. Sample Nd-doped BaTiO₃ ($0.3 \le x \le 0.10$) showed the conductivity of typical insulator [15] which at lower temperature, the conductivity lies in range of 10^{-8} to 10^{-7} Scm⁻¹. However, it shows the gradual increment with increasing temperatures. Therefore, total conductivity was affected by temperatures and frequencies.

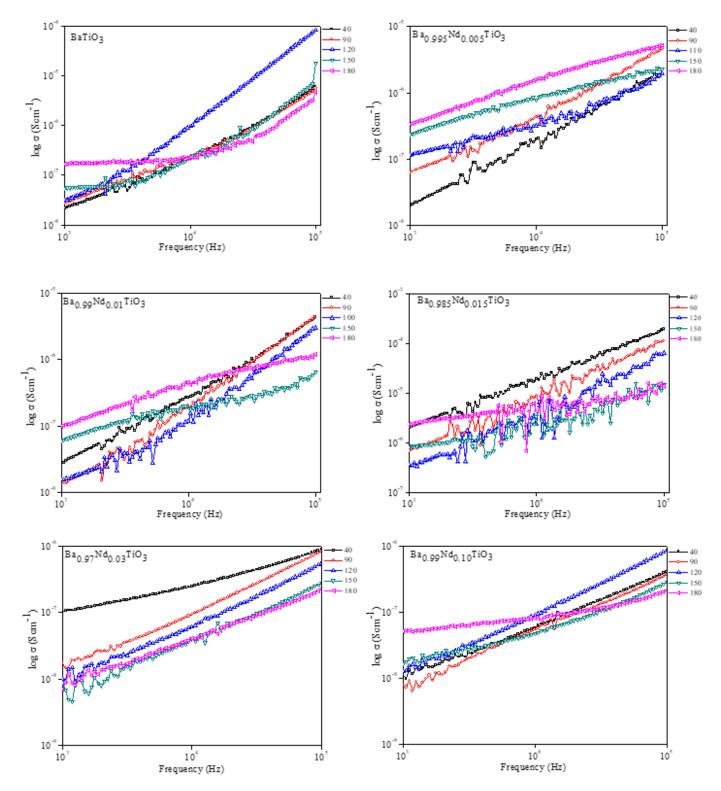


Figure 6. Conductivity of $Ba_{1-x}Nd_xTiO_3$ ($0 \le x \le 0.13$) at 1 kHz, 10 kHz and 100 kHz.

The piezoelectric charge constant (d_{33}) as a function of Nd addition is shown in Figure 7. With increasing Nd content (x), the d_{33} value decreases. However, for each sample, d_{33} increases with higher poling voltage. The reduction in d_{33} correlates with the decrease in permittivity upon Nd addition. The application of high poling voltage enhances the piezoelectric effect, likely due to improved domain alignment during the poling process in BaTiO₃ ceramics. It is presumed that the addition of Nd reduces the number of

domains that align during poling as reported by Ganguly et al [17]. Furthermore, the structural transition from the tetragonal to cubic phase weakens the ferroelectricity of the samples. Although high d_{33} is partly attributed to high permittivity, the Nd-doped BaTiO_3 samples (0.005 $\leq x \leq$ 0.015) exhibit lower d_{33} values. This may be due to the large strain observed in the grains of Nd-doped BaTiO_3, which also contributes to the reduced d_{33} values [17].

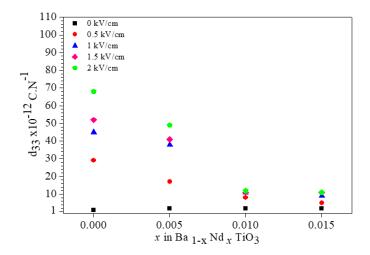


Figure 7. Piezoelectric charge constant, d_{33} of $Ba_{1-x}Nd_xTiO_3$ ($0 \le x \le 0.015$) in varied poling voltage.

Figure 8 shows the piezoelectric voltage constant (g₃₃) of Nd-doped BaTiO₃ as a function of poling voltage. At a maximum poling electric field of 2 kV/cm, the g₃₃ value increased significantly, reaching 7×10^{-2} V·m/N, 1×10^{-2} V·m/N, 0.7×10^{-2} V·m/N, and 0.5×10^{-2} V·m/N for various Nd concentrations. The piezoelectric voltage constant, g₃₃,

of both undoped and Nd-doped $BaTiO_3$ shows a notable decrease in parallel with the d_{33} values. This strongly suggests that both poling electric field and poling temperature significantly influence the g_{33} values, in addition to the intrinsic properties of the material.

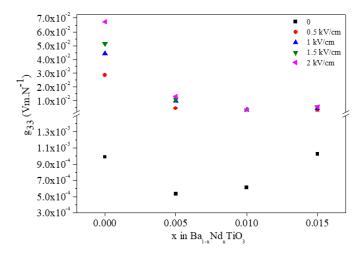


Figure 8. Piezoelectric voltage constant, g_{33} of $Ba_{1-x}Nd_xTiO_3$ ($0 \le x \le 0.015$) in varied poling voltage.

4. CONCLUSION

The modification of structural, electrical, and piezoelectric properties of neodymium-doped BaTiO₃ is attributed to the incorporation of Nd³⁺ ions into the Ba²⁺ sites. At low doping levels (small x), the tetragonal structure is retained in the XRD patterns, though accompanied by a reduction in lattice parameters and unit cell volume. Further increases in Nd content leads to the creation of Ba-site vacancies, which in turn initiate a structural transition from tetragonal to cubic. The solid solution range was determined to be $0 \le x \le 0.10$. Rietveld refinement was employed to model the [BaTiO₃], $[Nd_2O_3]$, and $[TiO_6]$ units within the crystal lattice. The refinement results showed relatively low χ^2 values and acceptable R-factors (below 20%). The Ti atom positions tend to shift toward the center of the octahedron due to the reduction in tetragonality associated with Nd^{3+} substitution for Ba²⁺.

Electrical characterization revealed that neodymium doping significantly lowered the Curie temperature (Tc) to below 30 °C. This shift in Tc is consistent with the observed decrease in permittivity as x increases. Neodymium also inhibits the formation of ferroelectric domains, suppressing the alignment of individual dipoles. Conductivity analysis showed that $\log \sigma$ exhibited both frequency and temperature dependence within the range $0 \le x \le 0.015$. The piezoelectric properties (such as d_{33}) also decreased significantly with increasing Nd content, likely due to reduced remanent polarization after the poling process. These findings strongly highlight the critical role of doping mechanisms in generating structural defects that alter the fundamental properties of pure BaTiO_3.

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